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The Teixeira singularity degeneracy and its bifurcation in PWL systems

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Introduction

- We consider 3D piecewise linear Filippov differential systems with a separation plane, having a two-fold point with invisible tangencies, that is, the so-called Teixeira singularity (TS-point, for short).
- For some parameter values this singularity undergoes a compound bifurcation: there appears a sliding bifurcation involving a pseudo-equilibrium point and, simultaneously, a bifurcation associated to the birth of a crossing limit cycle.
- After determining a generic canonical form, we show how to characterize such a compound bifurcation.
- Our motivation comes from the natural appearance of TS-points in the control of Boost converters.

Summary

- A non-trivial yet manageable single-parameter example
- The general DPWL case:
 - Hypotheses for having a TS-point and canonical form
 - The sliding dynamics bifurcation
 - The crossing dynamics bifurcation
- A multi-parametric example
- TS-point in a DC-DC Boost converter
- Conclusions





















The crossing dynamics bifurcation (periodic orbits?)

 \mathcal{Z}

 Σ_{as}

 Σ_c^-

 \boldsymbol{y}

 Σ_c^+

 Σ_{rs}

$$P_{+}\begin{pmatrix} y\\ z \end{pmatrix} = P_{-}^{-1} \begin{pmatrix} y\\ z \end{pmatrix}$$
$$\dot{x} = -y$$
$$\dot{y} = 1$$
$$\dot{z} = -2$$

The (+) vector field gives

$$\begin{aligned} x_{+}(t) \\ y_{+}(t) \\ z_{+}(t) \end{aligned} &= \begin{bmatrix} -t(t+2y_{0})/2 \\ t+y_{0} \\ -2t+z_{0} \end{bmatrix} \text{ with } \begin{bmatrix} x_{+}(0) \\ y_{+}(0) \\ z_{+}(0) \end{bmatrix} = \begin{bmatrix} 0 \\ y_{0} \\ z_{0} \end{bmatrix} \\ \text{ so that } P_{+} \begin{pmatrix} y_{0} \\ z_{0} \end{pmatrix} = \begin{pmatrix} -y_{0} \\ -4y_{0}+z_{0} \end{pmatrix} \end{aligned}$$

The crossing dynamics

The integration of the (-) vector field is much more involved but the choice of stepped eigenvalues allows to express solutions in an algebraic way.

$$A^{-} = \begin{pmatrix} -6a & 0 & 1\\ 6a^{3} & 0 & 0\\ -11a^{2} & -1 & 0 \end{pmatrix} \qquad \operatorname{Spec}(A^{-}) = \{-a, -2a, -3a\}$$

$$e^{-A^{-}t} = \frac{e^{at}}{2} \begin{pmatrix} 1 - 8e^{at} + 9e^{2at} & -\frac{(-1 + e^{at})^2}{a^2} & -\frac{(-1 + e^{at})(-1 + 3e^{at})}{a} \\ -6a^2(-1 + e^{at})(-1 + 3e^{at}) & 2\left(3 - 3e^{at} + e^{2at}\right) & 6a\left(-1 + e^{at}\right)^2 \\ a\left(-1 + e^{at}\right)(-5 + 27e^{at}) & -\frac{(-1 + e^{at})(-5 + 3e^{at})}{a} & -5 + 16e^{at} - 9e^{2at} \end{pmatrix} = \frac{1}{2} \left(\frac{1 - 8e^{at} + 9e^{2at}}{a} \right) = \frac{1}{2} \left(\frac{1 - 8e^{at} + 9e^{2at}}{a} + \frac{1 - 8e^{at}}{a} \right) = \frac{1 - 8e^{at} + 9e^{2at}}{a} = \frac{1 - 8e^{at} + 9e$$

$$=\frac{1}{2u^3}\left(\begin{array}{ccc}u^2-8u+9&-\frac{(u-1)^2}{a^2}&-\frac{(u-3)(u-1)}{a}\\-6a^2(u-3)(u-1)&2\left(3u^2-3u+1\right)&6a(u-1)^2\\a(u-1)(5u-27)&-\frac{(u-1)(5u-3)}{a}&-5u^2+16u-9\end{array}\right)$$

$$0 < u = e^{-at} < 1, \quad a > 0$$

Looking for $P_{-}^{-1}(y, z)$ we write

$$\begin{bmatrix} x_{-}^{-1}(t) \\ y_{-}^{-1}(t) \\ z_{-}^{-1}(t) \end{bmatrix} = \frac{1}{6a^2} \begin{bmatrix} 1 \\ -5a^2 \\ 6a \end{bmatrix} + e^{-A^-t} \left(\begin{bmatrix} 0 \\ y_0 \\ z_0 \end{bmatrix} - \frac{1}{6a^2} \begin{bmatrix} 1 \\ -5a^2 \\ 6a \end{bmatrix} \right)$$

The condition $x_{-}^{-1}(t) = 0$ gives $0 < u = e^{-at} < 1, a > 0$

$$a = \frac{(-1+u)(2+u-3y_0)}{3(-3+u)z_0}, \text{ and then}$$
$$y_{-}^{-1}(u) = \frac{-1-9u+9u^2+u^3+6(1-3u)y_0}{6(-3+u)u^2},$$
$$z_{-}^{-1}(u) = \frac{(-1-6u+u^2+6y_0)z_0}{2u^2(2+u-3y_0)}$$

Finally, imposing

$$P_{+}\begin{pmatrix}y_{0}\\z_{0}\end{pmatrix} = \begin{bmatrix}-2y_{0}\\-4y_{0}+z_{0}\end{bmatrix} = P_{-}^{-1}\begin{pmatrix}y_{0}\\z_{0}\end{pmatrix} = \begin{bmatrix}y_{-}^{-1}(u)\\z_{-}^{-1}(u)\end{bmatrix}$$

we get

$$a = \frac{1 + 2u + 6u^2 + 2u^3 + u^4}{2u(1 + 10u + u^2)}, \text{ and then}$$
$$y_0 = -\frac{(-1 + u)(1 + 10u + u^2)}{6(1 + u)(1 - 4u + u^2)},$$
$$z_0 = \frac{(-1 + u)(-1 + 6u + 5u^2 + 2u^3)}{6a(1 + u)(1 - 4u + u^2)}$$

Note that $1 - 4u + u^2 = (u - 2)^2 - 3$, so that the above expressions are only valid for $2 - \sqrt{3} < u < 1$. Also, $\lim_{u \to 1^-} a = 1/2$, $\lim_{u \to (2 - \sqrt{3})^+} a = \infty$.

A periodic orbit exists for 0.5 < a < 1, corresponding to a decreasing of u in the interval $2 - \sqrt{3} < u < 1$. The periodic orbit is born from the TS-point at a =0.5 and disappears in a bifurcation at infinity for a = 1.



A periodic orbit exists for 0.5 < a < 1, 1.0 corresponding to a decreasing of u in the interval $2 - \sqrt{3} < u < 1$. The periodic orbit is born from the TS-point at a =0.5 and disappears in a bifurcation at 0.8 infinity for a = 1. \mathcal{U} 0.6 We can compute in terms of u, the determinant and trace of the derivative $DP(y_0, z_0) = DP_{-}^{-1} \circ DP_{+}(y_0, z_0)$ to 0.4 check the stability of the periodic orbit. $2 - \sqrt{3}$ 0.7 0.5 0.6 0.8 0.9 1.0 \boldsymbol{a}

The stability triangle in the (trace,det)-plane



The followed path for 0.5< *a* <1 in the (trace,det)-plane



The red path in the (trace,det)-plane starts at the Bogdanov-Takens point (2,1). Note that the final point in the path is over the horizontal axis, namely

$$(98 - 56\sqrt{3}, 97 - 56\sqrt{3}).$$

For this example, the crossing periodic orbit which is born from the TS-point is always of **stable node** type. Such a global control of the compound bifurcation for the Teixeira singularity is rather unusual: in what follows, we show how it is possible to characterize locally this bifurcation for discontinuous linear systems.

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The general DPWL case

There exist three vectors \mathbf{w}^- , \mathbf{w}^+ , $\mathbf{v} \in \mathbb{R}^3$, with $\mathbf{v} \neq \mathbf{0}$, two 3×3 square matrices A^- and A^+ and a scalar $\delta \in \mathbb{R}$, such that

(a)
$$\mathbf{F}(\mathbf{x}) = \begin{cases} \mathbf{F}^{-}(\mathbf{x}) = A^{-}\mathbf{x} + \mathbf{w}^{-}, & \text{if } \mathbf{v}^{T}\mathbf{x} + \delta < 0 \\ \mathbf{F}^{+}(\mathbf{x}) = A^{+}\mathbf{x} + \mathbf{w}^{+}, & \text{if } \mathbf{v}^{T}\mathbf{x} + \delta > 0 \end{cases}$$

(b) $A^{-}\mathbf{x} + \mathbf{w}^{-} \neq A^{+}\mathbf{x} + \mathbf{w}^{+}$ generically, when $\mathbf{v}^{T}\mathbf{x} + \delta = 0$.

 $\dot{\mathbf{x}} = \mathbf{F}^{-}(\mathbf{x})$

$$\dot{\mathbf{x}} = \mathbf{F}^+(\mathbf{x})$$

The switching manifold is the plane

$$\Sigma = \{ \mathbf{x} \in \mathbb{R}^3 : \mathbf{v}^T \mathbf{x} + \delta = 0 \}$$

Hypotheses for having a TS-point

(H1) The tangency lines

$$T^{-} = \{ \mathbf{x} \in \Sigma : \langle \mathbf{v}, \mathbf{F}^{-}(\mathbf{x}) \rangle = 0 \}$$
$$T^{+} = \{ \mathbf{x} \in \Sigma : \langle \mathbf{v}, \mathbf{F}^{+}(\mathbf{x}) \rangle = 0 \}$$

intersect transversally at a point $\hat{\mathbf{x}}$.

(H2) At the two-fold point $\hat{\mathbf{x}}$ both tangencies are invisible: $\mathbf{v}^T A^- \mathbf{F}^-(\hat{\mathbf{x}}) > 0$ and $\mathbf{v}^T A^+ \mathbf{F}^+(\hat{\mathbf{x}}) < 0$.

The canonical form for DPWLS with a TS-point

Proposition. Under hypotheses (H1) and (H2), it is possible through an invertible linear change of coordinates and a rescaling of time, to rewrite the system in the canonical form

$$\dot{\mathbf{x}} = \begin{cases} \mathbf{F}^{-}(\mathbf{x}), & \text{if } x < 0, \\ \mathbf{F}^{+}(\mathbf{x}), & \text{if } x > 0, \end{cases}$$
(1)

with $\mathbf{x} = (x, y, z) \in \mathbb{R}^3$, where the linear vector fields $\mathbf{F}^{\pm} : \mathbb{R}^3 \longrightarrow \mathbb{R}^3$ are

$$\mathbf{F}^{+}(\mathbf{x}) = \begin{bmatrix} f_{1}^{+} & -1 & 0\\ g_{1}^{+} & g_{2}^{+} & g_{3}^{+}\\ h_{1}^{+} & h_{2}^{+} & h_{3}^{+} \end{bmatrix} \begin{bmatrix} x\\ y\\ z \end{bmatrix} + \begin{bmatrix} 0\\ 1\\ v_{+} \end{bmatrix}, \quad \mathbf{F}^{-}(\mathbf{x}) = \begin{bmatrix} f_{1}^{-} & 0 & 1\\ g_{1}^{-} & g_{2}^{-} & g_{3}^{-}\\ h_{1}^{-} & h_{2}^{-} & h_{3}^{-} \end{bmatrix} \begin{bmatrix} x\\ y\\ z \end{bmatrix} + \begin{bmatrix} 0\\ v_{-}\\ 1 \end{bmatrix},$$

for some constants f_1^{\pm} , g_1^{\pm} , g_2^{\pm} , g_3^{\pm} , h_1^{\pm} , h_2^{\pm} , h_3^{\pm} and v_{\pm} .

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for some constants f_1^{\pm} , g_1^{\pm} , g_2^{\pm} , g_3^{\pm} , h_1^{\pm} , h_2^{\pm} , h_3^{\pm} and v_{\pm} .

 $T_{+} = \{(x, y, z) \in \mathbb{R}^{3} : x = y = 0\} \qquad T_{-} = \{(x, y, z) \in \mathbb{R}^{3} : x = z = 0\}$ $\Sigma = \{\mathbf{x} = (x, y, z) \in \mathbb{R}^{3} : x = 0\} \qquad \text{TS-point at the origin!}$

The canonical form for DPWLS with a TS-point

$$\Sigma_{as} = \{(x, y, z) \in \mathbb{R}^3 : x = 0, y > 0 \text{ and } z > 0\}$$

$$\Sigma_{rs} = \{(x, y, z) \in \mathbb{R}^3 : x = 0, y < 0 \text{ and } z < 0\}$$

$$\Sigma_c^- = \{(x, y, z) \in \mathbb{R}^3 : x = 0, y > 0 \text{ and } z < 0\}$$

$$\Sigma_c^+ = \{(x, y, z) \in \mathbb{R}^3 : x = 0, y < 0 \text{ and } z > 0\}$$



The sliding vector field associated to system (1) is

$$\mathbf{F}^{s}(0,y,z) = \frac{1}{y+z} \begin{bmatrix} 0 \\ v_{-}y+z+g_{2}^{-}y^{2}+(g_{2}^{+}+g_{3}^{-})yz+g_{3}^{+}z^{2} \\ y+v_{+}z+h_{2}^{-}y^{2}+(h_{2}^{+}+h_{3}^{-})yz+h_{3}^{+}z^{2} \end{bmatrix}$$

We work with the desingularized system

$$\dot{\mathbf{x}} = \mathbf{F}_{d}^{s}(0, y, z) := \begin{bmatrix} 0 \\ v_{-}y + z + g_{2}^{-}y^{2} + (g_{2}^{+} + g_{3}^{-})yz + g_{3}^{+}z^{2} \\ y + v_{+}z + h_{2}^{-}y^{2} + (h_{2}^{+} + h_{3}^{-})yz + h_{3}^{+}z^{2} \end{bmatrix}$$
(2)

Pseudo-equilibrium points come from solving the two equations

$$v_{-}y + z + \begin{bmatrix} y & z \end{bmatrix} G \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$
$$y + v_{+}z + \begin{bmatrix} y & z \end{bmatrix} H \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$

where

$$G = \begin{bmatrix} g_2^- & g_3^- \\ g_2^+ & g_3^+ \end{bmatrix} \text{ and } H = \begin{bmatrix} h_2^- & h_3^- \\ h_2^+ & h_3^+ \end{bmatrix}$$

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Pseudo-equilibrium points come from solving the two equations

$$J(0,0) = \begin{bmatrix} v_{-} & 1 \\ 1 & v_{+} \end{bmatrix} \leftarrow \begin{bmatrix} v_{-}y + z + \begin{bmatrix} y & z \end{bmatrix} G \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$
$$y + v_{+}z + \begin{bmatrix} y & z \end{bmatrix} H \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$

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Apart from the origin, we can have another non-trivial emanating branch for $v_{-}v_{+} = 1$ (transcritical bifurcation). We assume all parameters fixed excepting

$$v_{-} = v_{-}(\varepsilon) := \frac{1+\varepsilon}{v_{+}}$$

Pseudo-equilibrium points come from solving the two equations

$$J(0,0) = \begin{bmatrix} v_{-} & 1 \\ 1 & v_{+} \end{bmatrix} \leftarrow \begin{bmatrix} v_{-}y + z + \begin{bmatrix} y & z \end{bmatrix} G \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$
$$y + v_{+}z + \begin{bmatrix} y & z \end{bmatrix} H \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$

where

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$$\underbrace{(v_+v_--1)}_{\varepsilon} y + \begin{bmatrix} y & z \end{bmatrix} (v_+G - H) \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$
$$y + v_+z + \begin{bmatrix} y & z \end{bmatrix} H \begin{bmatrix} y \\ z \end{bmatrix} = 0.$$

We can use the equivalent system of equations

$$\varepsilon y + \begin{bmatrix} y & z \end{bmatrix} (v_{+}G - H) \begin{bmatrix} y \\ z \end{bmatrix} = 0,$$
$$y + v_{+}z + \begin{bmatrix} y & z \end{bmatrix} H \begin{bmatrix} y \\ z \end{bmatrix} = 0.$$

Using the implicit function theorem for the second equation at (y, z) = (0, 0)we get that, for z small, solutions must satisfy

$$\begin{bmatrix} y \\ z \end{bmatrix} = \begin{bmatrix} -v_+ \\ 1 \end{bmatrix} z + O(z^2).$$
 We need $v_+ < 0$

Substituting this expansion in the first equation, desingularizing it and assuming

$$\kappa_S = \begin{bmatrix} -v_+ & 1 \end{bmatrix} (v_+ G - H) \begin{bmatrix} -v_+ \\ 1 \end{bmatrix} \neq 0,$$

we get

$$z = \frac{v_+}{\kappa_S}\varepsilon + O(\varepsilon^2).$$

Proposition (Pseudo-equilibrium transition) Assuming that the criticality coefficient

$$\kappa_S = \begin{bmatrix} -v_+ & 1 \end{bmatrix} (v_+ G - H) \begin{bmatrix} -v_+ \\ 1 \end{bmatrix} \neq 0$$

the following statements hold.

- (a) System (2) undergoes a transcritical bifurcation for $\varepsilon = 0$, so that there exists a branch of equilibria $(\tilde{y}(\varepsilon), \tilde{z}(\varepsilon))$ with $(\tilde{y}(0), \tilde{z}(0)) = (0, 0)$ and $(\tilde{y}'(0), \tilde{z}'(0)) = \left(\frac{-v_+^2}{\kappa_S}, \frac{v_+}{\kappa_S}\right)$.
- (b) For the particular case where $v_+ < 0$, the emanating branch is located at the quadrants with yz > 0. If $\kappa_S > 0$ ($\kappa_S < 0$) then in passing from $\varepsilon < 0$ to $\varepsilon > 0$ the origin passes from being a saddle to a stable node, while the nontrivial equilibrium passes from being a stable node in the first (third) quadrant to be a saddle in the third (first) quadrant.

Theorem S Assuming $v_+ < 0$, $v_- < 0$ and $\kappa_S \neq 0$, system (1) has for $\varepsilon = v_-v_+ - 1$ with $|\varepsilon| > 0$ small, one pseudo-equilibrium point $\tilde{\mathbf{x}}(\varepsilon) = (0, \tilde{y}(\varepsilon), \tilde{z}(\varepsilon))$, such that

$$(\tilde{y}(\varepsilon), \tilde{z}(\varepsilon)) = \left(-\frac{v_+^2}{\kappa_S}, \frac{v_+}{\kappa_S}\right)\varepsilon + O(\varepsilon^2),$$

and the following statements hold.

- (a) (Supercritical case) If $\kappa_S > 0$, then $\widetilde{\mathbf{x}}(\varepsilon) \in \Sigma_{as}$ is a stable pseudo-node for $\varepsilon < 0$, being $\widetilde{\mathbf{x}}(\varepsilon) \in \Sigma_{rs}$ a pseudo-saddle for $\varepsilon > 0$.
- (b) (Subcritical case) If $\kappa_S < 0$, then $\widetilde{\mathbf{x}}(\varepsilon) \in \Sigma_{rs}$ is an unstable pseudo-node for $\varepsilon < 0$, being $\widetilde{\mathbf{x}}(\varepsilon) \in \Sigma_{as}$ a pseudo-saddle for $\varepsilon > 0$.


The crossing dynamics bifurcation (periodic orbits?)

$$P_+\begin{pmatrix}y\\z\end{pmatrix} = P_-^{-1}\begin{pmatrix}y\\z\end{pmatrix}$$



The crossing dynamics bifurcation (periodic orbits?)

$$P_{+}\begin{pmatrix} y\\z \end{pmatrix} = P_{-}^{-1}\begin{pmatrix} y\\z \end{pmatrix} = P_{-}\begin{pmatrix} y\\z \end{pmatrix}$$
Involution property



The crossing dynamics bifurcation (periodic orbits?)

$$P_{+}\begin{pmatrix} y\\ z \end{pmatrix} = P_{-}^{-1}\begin{pmatrix} y\\ z \end{pmatrix} = P_{-}\begin{pmatrix} y\\ z \end{pmatrix}$$
Involution property
$$P_{+}\begin{pmatrix} \begin{bmatrix} 0\\ z \end{bmatrix} \end{pmatrix} = \begin{bmatrix} 0\\ z \end{bmatrix}; \quad P_{-}\begin{pmatrix} \begin{bmatrix} y\\ 0 \end{bmatrix} \end{pmatrix} = \begin{bmatrix} y\\ 0 \end{bmatrix}$$
Invariance of the tangency lines



The structure of return maps

$$P_{+}\begin{pmatrix}y\\z\end{pmatrix} = \begin{bmatrix}-1 & 0\\-2v_{+} & 1\end{bmatrix}\begin{bmatrix}y\\z\end{bmatrix} + y\begin{bmatrix}q_{11}^{+}y + q_{12}^{+}z\\q_{21}^{+}y + q_{22}^{+}z\end{bmatrix} + y\begin{bmatrix}c_{11}^{+}y^{2} + c_{12}^{+}yz + c_{13}^{+}z^{2}\\c_{21}^{+}y^{2} + c_{22}^{+}yz + c_{23}^{+}z^{2}\end{bmatrix} + O(4)$$

$$P_{-}\begin{pmatrix}y\\z\end{pmatrix} = \begin{bmatrix}1 & -2v_{-}\\0 & -1\end{bmatrix}\begin{bmatrix}y\\z\end{bmatrix} + z\begin{bmatrix}q_{11}^{-}y + q_{12}^{-}z\\q_{21}^{-}y + q_{22}^{-}z\end{bmatrix} + z\begin{bmatrix}c_{11}^{-}y^{2} + c_{12}^{-}yz + c_{13}^{-}z^{2}\\c_{21}^{-}y^{2} + c_{22}^{-}yz + c_{23}^{-}z^{2}\end{bmatrix} + O(4)$$

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$$\begin{bmatrix}q_{12}^{+} = 0;\\q_{21}^{+} = v_{+}(q_{11}^{+} - q_{22}^{+});\\c_{11}^{+} = -v_{+}c_{12}^{+} - (q_{11}^{+})^{2};\\c_{13}^{+} = 0;\\c_{22}^{+} = v_{+}(c_{12}^{+} - 2c_{23}^{+}) + q_{22}^{+}(q_{22}^{+} - q_{11}^{+})/2.$$

$$Involution property$$

$$P_{-}\begin{pmatrix}y\\z\end{pmatrix} = \begin{bmatrix}1 & -2v_{-}\\0 & -1\end{bmatrix}\begin{bmatrix}y\\z\end{bmatrix} + z\begin{bmatrix}q_{11}^{-}y + q_{12}^{-}z\\q_{21}^{-}y + q_{22}^{-}z\end{bmatrix} + z\begin{bmatrix}c_{11}^{-}y^{2} + c_{12}^{-}yz + c_{13}^{-}z^{2}\\c_{21}^{-}y^{2} + c_{22}^{-}yz + c_{23}^{-}z^{2}\end{bmatrix} + O(4)$$

The structure of return maps

$$P_{+}\begin{pmatrix}y\\z\end{pmatrix} = \begin{bmatrix}-1 & 0\\-2v_{+} & 1\end{bmatrix}\begin{bmatrix}y\\z\end{bmatrix} + y\begin{bmatrix}q_{11}^{+}y + \dot{q}_{12}z\\\dot{q}_{21}^{+}y + \dot{q}_{22}^{+}z\end{bmatrix} + y\begin{bmatrix}\dot{q}_{11}^{+}y^{2} + c_{12}^{+}yz + \dot{q}_{32}z^{2}\\c_{21}^{+}y^{2} + \dot{q}_{22}^{+}yz + c_{23}^{+}z^{2}\end{bmatrix} + O(4)$$

$$\begin{bmatrix}q_{12}^{+} = 0;\\q_{21}^{+} = v_{+}(q_{11}^{+} - q_{22}^{+});\\c_{11}^{+} = -v_{+}c_{12}^{+} - (q_{11}^{+})^{2};\\c_{13}^{+} = 0;\\c_{22}^{+} = v_{+}(c_{12}^{+} - 2c_{23}^{+}) + q_{22}^{+}(q_{22}^{+} - q_{11}^{+})/2. \end{bmatrix}$$
Involution property
$$P_{-}\begin{pmatrix}y\\z\end{pmatrix} = \begin{bmatrix}1 & -2v_{-}\\0 & -1\end{bmatrix}\begin{bmatrix}y\\z\end{bmatrix} + z\begin{bmatrix}q_{11}^{-}y + \dot{q}_{12}yz\\\dot{q}y + q_{22}z\end{bmatrix} + z\begin{bmatrix}c_{11}^{-}y^{2} + \dot{q}_{12}yz + c_{13}z^{2}\\\dot{q}y^{2} + c_{22}yz + \dot{q}_{23}z^{2}\end{bmatrix} + O(4)$$

$$\begin{bmatrix}q_{21}^{-} = 0;\\q_{12}^{-} = v_{-}(q_{22}^{-} - q_{11}^{-});\\c_{23}^{-} = -v_{-}c_{22}^{-} - (q_{22}^{-})^{2};\\c_{21}^{-} = 0;\\c_{12}^{-} = v_{-}(c_{22}^{-} - 2c_{11}^{-}) + q_{11}^{-}(q_{11}^{-} - q_{22}^{-})/2. \end{bmatrix}$$
Involution property

Looking for non-trivial fixed points

We impose the equality $P_+(y,z) = P_-^{-1}(y,z)$, to obtain

$$0 = -2y + 2v_{-}z + q_{11}^{+}y^{2} - q_{11}^{-}yz - q_{12}^{-}z^{2} + c_{11}^{+}y^{3} + (c_{12}^{+} - c_{11}^{-})y^{2}z - c_{12}^{-}yz^{2} - c_{13}^{-}z^{3} + O(4),$$

$$0 = -2v_{+}y + 2z + q_{21}^{+}y^{2} + q_{22}^{+}yz - q_{22}^{-}z^{2} + c_{21}^{+}y^{3} + c_{22}^{+}y^{2}z + (c_{23}^{+} - c_{22}^{-})yz^{2} - c_{23}^{-}z^{3} + O(4).$$

Again, we can have a bifurcation at $v_+v_- = 1$ so that we do a new bifurcation analysis by assuming all parameters fixed, excepting v_- , and take

$$v_{-} = \frac{1+\varepsilon}{v_{+}}$$

Note that, from the second equation and the implicit function theorem, we can assume the existence of a function $p(y, \varepsilon)$ such that

 $z = y \cdot p(y, \varepsilon)$, with $p(0, 0) = v_+$.

We need
$$v_+ < 0$$

Looking for non-trivial fixed points

We get
$$z = y \cdot p(y, \varepsilon)$$
 with

$$p(y, \varepsilon) = v_{+} - \left(q_{11}^{+} - q_{22}^{-}(\varepsilon)v_{+}\right)\frac{y}{2} + \left[-2c_{21}^{+} + q_{22}^{+}(2q_{11}^{+} - q_{22}^{+})v_{+} + \left(2c_{22}^{-}(\varepsilon) - 2c_{12}^{+} + 2c_{23}^{+} - 2q_{22}^{-}(\varepsilon)q_{11}^{+} - q_{22}^{-}(\varepsilon)q_{22}^{+}\right)v_{+}^{2} - 2c_{22}^{-}(\varepsilon)v_{-}(\varepsilon)v_{+}^{3}\right]\frac{y^{2}}{4} + \dots$$

After substituting z in the first equation, we can desingularize it to get

$$0 = 2\varepsilon - (q_{11}^+ - q_{11}^-(\varepsilon)v_+)\varepsilon y + P(\varepsilon)y^2 + \dots$$

with

$$\begin{split} P(\varepsilon) &= -\frac{1}{2v_{+}} \left[2c_{21}^{+} + \left(2(q_{11}^{+})^{2} - 2q_{11}^{+}q_{22}^{+} + (q_{22}^{+})^{2} \right) v_{+} + \right. \\ &+ \left(-2c_{11}^{-}(\varepsilon) + 2c_{22}^{-}(\varepsilon) + 2c_{12}^{+} - 2c_{23}^{+} + q_{11}^{-}(\varepsilon)q_{11}^{+} + q_{22}^{-}(\varepsilon)q_{22}^{+} \right) v_{+}^{2} + \\ &+ \left(q_{11}^{-}(\varepsilon)^{2} - 2q_{11}^{-}(\varepsilon)q_{22}^{-}(\varepsilon) + 2q_{22}^{-}(\varepsilon)^{2} \right) v_{+}^{3} + 2c_{13}^{-}(\varepsilon)v_{+}^{4} + O(\varepsilon) \right] \end{split}$$

Looking for non-trivial fixed points

Clearly, it is more convenient to parameterize the emanating branch in terms of y, so that we get the expansion

$$\varepsilon = \kappa_C \cdot y^2 + O(y^3),$$

where the criticality coefficient is

$$\kappa_{C} = \frac{1}{4v_{+}} \left[2c_{21}^{+} + \left(2(q_{11}^{+})^{2} - 2q_{11}^{+}q_{22}^{+} + (q_{22}^{+})^{2} \right) v_{+} + \left(-2c_{11}^{-}(0) + 2c_{22}^{-}(0) + 2c_{12}^{+} - 2c_{23}^{+} + q_{11}^{-}(0)q_{11}^{+} + q_{22}^{-}(0)q_{22}^{+} \right) v_{+}^{2} + \left(q_{11}^{-}(0)^{2} - 2q_{11}^{-}(0)q_{22}^{-}(0) + 2q_{22}^{-}(0)^{2} \right) v_{+}^{3} + 2c_{13}^{-}(0)v_{+}^{4} \right]$$

In short, depending on the sign of κ_C we have a subcritical or supercritical bifurcation of non-trivial fixed points with

$$z(y) = v_+ \cdot y + O(y^2),$$

$$\varepsilon(y) = \kappa_C \cdot y^2 + O(y^3).$$



Topological type of non-trivial fixed points

Just computing the derivatives of the Poincaré half-return maps, and evaluating them at the branch of non-trivial fixed points, we get $DP(y) = D(P_{-} \circ P_{+})(y, z(y))$ and the expansions for its determinant ant trace, namely

det
$$DP(y) = 1 + d_1y + d_2y^2 + O(y^3)$$
,
trace $DP(y) = 2 + t_1y + t_2y^2 + O(y^3)$.

It turns out that $t_1 = d_1$. When this common value vanishes, we have a degeneracy that should require much more long computations. We define $\sigma := t_1 = d_1$, and the computations give

$$\begin{split} \sigma &= -2q_{11}^{+} + q_{22}^{+} + \left(2q_{22}^{-}(0) - q_{11}^{-}(0)\right)v_{+}, \\ d_{2} &= \frac{1}{2}\left[6(q_{11}^{+})^{2} - 5q_{11}^{+}\left(q_{22}^{+} - q_{11}^{-}(0)v_{+} + 2q_{22}^{-}(0)v_{+}\right) + \left(q_{22}^{+} - q_{11}^{-}(0)v_{+} + 2q_{22}^{-}(0)v_{+}\right)^{2}\right], \\ t_{2} &= -\frac{1}{2v_{+}}\left[8c_{21}^{+} + \left(2(q_{11}^{+})^{2} - 3q_{11}^{+}q_{22}^{+} + 3(q_{22}^{+})^{2}\right)v_{+} + \right. \\ &+ \left(-8c_{11}^{-}(0) + 8c_{22}^{-}(0) + 8c_{12}^{+} - 8c_{23}^{+} - q_{11}^{-}(0)q_{11}^{+} + 10q_{22}^{-}(0)q_{11}^{+} + 2q_{11}^{-}(0)q_{22}^{+}\right)v_{+}^{2} + \\ &+ \left(3q_{11}^{-}(0)^{2} - 4q_{11}^{-}(0)q_{22}^{-}(0) + 4q_{22}^{-}(0)^{2}\right)v_{+}^{3} + 8c_{13}^{-}(0)v_{+}^{4}\right] \end{split}$$

Topological type of non-trivial fixed points We have at the branch of non-trivial fixed points det $DP(u) = 1 + \sigma u + d_0 u^2 + O(u^3)$

trace
$$DP(y) = 1 + \sigma y + u_2 y^2 + O(y^3)$$
,
 $trace DP(y) = 2 + \sigma y + t_2 y^2 + O(y^3)$.

and, surprisingly, we get that $d_2 - t_2 = 8\kappa_C$.

Topological type of non-trivial fixed points We have at the branch of non-trivial fixed points

det
$$DP(y) = 1 + \sigma y + d_2 y^2 + O(y^3)$$
,
trace $DP(y) = 2 + \sigma y + t_2 y^2 + O(y^3)$.

and, surprisingly, we get that $d_2 - t_2 = 8\kappa_C$.



Theorem C Assume in system (1) that $v_+ < 0$, $v_- < 0$ with $v_-v_+-1 = 0$ and that the two conditions $\sigma \neq 0$ and $\kappa_S \neq 0$ hold. By moving v_- and using the bifurcation parameter $\varepsilon = v_-v_+ - 1$, one crossing periodic orbit bifurcates from the origin for $\kappa_S \cdot \varepsilon > 0$ small.

The bifurcating periodic orbit is stable whenever $\sigma > 0$ and $\kappa_S > 0$.

The topological type of the corresponding fixed point for the Poincaré map is saddle, node or focus depending on whether $\kappa_S < 0$, $0 < \kappa_S < \sigma^2/32$ or $\kappa_S > \sigma^2/32$, respectively.



Computing the return maps

To compute the return map P_+ we write

$$\mathbf{x}(\tau) = e^{A^+\tau} \mathbf{x}_0 + \int_0^\tau e^{A^+(\tau-s)} \mathbf{v}^+ ds,$$

where

$$A^{+} = \begin{bmatrix} f_{x}^{+} & -1 & 0\\ g_{x}^{+} & g_{y}^{+} & g_{z}^{+}\\ h_{x}^{+} & h_{y}^{+} & h_{z}^{+} \end{bmatrix}, \quad \mathbf{x}_{0} = \begin{bmatrix} 0\\ y_{0}\\ z_{0} \end{bmatrix}, \quad \mathbf{v}^{+} = \begin{bmatrix} 0\\ 1\\ v_{+} \end{bmatrix},$$

being $y_0 < 0$. In practice, we take

$$\mathbf{x}(\tau) = \mathbf{x}_0 + \left(\tau I + \frac{\tau^2}{2}A^+ + \frac{\tau^3}{6}(A^+)^2 + \frac{\tau^4}{24}(A^+)^3\right)M + \mathcal{O}(\tau^5)$$

where $M = A^+ \mathbf{x}_0 + \mathbf{v}^+$ and I is the identity matrix of order 3.

From the first component we can determine an expression for the time $\tau_+ = \tau_+(y_0, z_0)$ such that $x(\tau_+) = 0$. The third order polynomial approximation is given by

$$\tau_+(y_0, z_0) = a_{10}y_0 + a_{20}y_0^2 + a_{11}y_0z_0 + a_{30}y_0^3 + a_{21}y_0^2z_0 + a_{12}y_0z_0^2,$$

with

$$\begin{split} a_{10} &= -2, \\ a_{20} &= \frac{2}{3}(f_1^+ + g_2^+ - 2g_3^+ v_+), \\ a_{11} &= 2g_3^+, \\ a_{20} &= \frac{2}{3}(f_1^+ + g_2^+ - 2g_3^+ v_+), \\ a_{30} &= -\frac{2}{9}\left(2(f_1^+ + g_2^+)^2 - 3(f_1^+ g_2^+ - g_3^+ h_2^+ + g_1^+) - \right. \\ &\left. -(5f_1^+ + 5g_2^+ + 3h_3^+)g_3^+ v_+ + 8(g_3^+ v_+)^2\right), \\ a_{12} &= -2(g_3^+)^2, \\ a_{21} &= \frac{4}{3}g_3^+ (3g_3^+ v_+ - f_1^+ - g_2^+ - h_3^+). \end{split}$$

The half-return map $(y_1, z_1) = P_+(y_0, z_0)$ satisfies

$$y_{1} = -y_{0} + q_{11}^{+} y_{0}^{2} + y_{0} (c_{11}^{+} y_{0}^{2} + c_{12}^{+} y_{0} z_{0}) + O(4),$$

$$z_{1} = -2v_{+} y_{0} + z_{0} + y_{0} (q_{21}^{+} y_{0} + q_{22}^{+} z_{0}) + y_{0} (c_{21}^{+} y_{0}^{2} + c_{22}^{+} y_{0} z_{0} + c_{23}^{+} z_{0}^{2}) + O(4),$$

where

$$\Rightarrow q_{11}^{+} = \frac{2}{3} \left(f_1^{+} + g_2^{+} + g_3^{+} v_+ \right),$$

$$\Rightarrow q_{22}^{+} = -2(h_3^{+} - g_3^{+} v_+),$$

$$q_{21}^{+} = v_+(q_{11}^{+} - q_{22}^{+}) = \frac{2}{3} \left(f_1^{+} + g_2^{+} - 2g_3^{+} v_+ + 3h_3^{+} \right) v_+,$$

$$c_{11}^{+} = -v_+ c_{12}^{+} - (q_{11}^{+})^2 = -\frac{2}{9} \left[2(f_1^{+} + g_2^{+})^2 + (f_1^{+} + g_2^{+} + 3h_3^{+})g_3^{+} v_+ - 4(g_3^{+} v_+)^2 \right],$$

$$\Rightarrow c_{12}^{+} = -\frac{2}{3} \left(f_1^{+} + g_2^{+} + 2g_3^{+} v_+ - h_3^{+} \right) g_3^{+},$$

$$\Rightarrow c_{21}^{+} = -\frac{2}{9} \left\{ 3(h_1^{+} + f_1^{+} h_2^{+} - h_2^{+} h_3^{+}) + \left[2(f_1^{+} + g_2^{+})^2 - 3(f_1^{+} g_2^{+} + g_1^{+} - g_3^{+} h_2^{+}) + \right. \\ \left. + 6h_3^{+}(f_1^{+} + g_2^{+} + h_3^{+}) \right] v_+ + 5g_3^{+}(f_1^{+} + g_2^{+} + 3h_3^{+})v_+^2 + 8(g_3^{+})^2 v_+^3 \right\},$$

$$c_{22}^{+} = v_+(c_{12}^{+} - 2c_{23}^{+}) + q_{22}^{+}(q_{22}^{+} - q_{11}^{+})/2 =$$

$$= \frac{2}{3} \left[h_3^{+}(f_1^{+} + g_2^{+} + 3h_3^{+}) - 2g_3^{+}(f_1^{+} + g_2^{+} + 5h_3^{+})v_+ \right] + 4(g_3^{+} v_+)^2,$$

$$\Rightarrow c_{23}^{+} = 2g_3^{+}(h_3^{+} - g_3^{+} v_+) = -g_3^{+}q_{22}^{+}.$$

Analogously, we can determine the half-return map $(y_2, z_2) = P_{-}^{-1}(y_0, z_0)$, getting

$$y_{2} = y_{0} - 2v_{-}z_{0} + z_{0}(q_{11}^{-}y_{0} + q_{12}^{-}z_{0}) + z_{0}(c_{11}^{-}y_{0}^{2} + c_{12}^{-}y_{0}z_{0} + c_{13}^{-}z_{0}^{2}) + O(4)$$

$$z_{2} = -z_{0} + q_{22}^{-}z_{0}^{2} + z_{0}(c_{22}^{-}y_{0}z_{0} + c_{23}^{-}z_{0}^{2}) + O(4),$$

where

$$\begin{array}{l} \bigstar \quad q_{11}^- = -2(g_2^- - h_2^- v_-), \\ q_{12}^- = v_-(q_{22}^- - q_{11}^-) = \frac{2}{3} \left(f_1^- + h_3^- - 2h_2^- v_- + 3g_2^- \right) v_-, \\ \clubsuit \quad q_{22}^- = \frac{2}{3} \left(f_1^- + h_3^- + h_2^- v_- \right), \\ \clubsuit \quad c_{11}^- = 2h_2^- (g_2^- - h_2^- v_-) = -h_2^- q_{11}^-, \\ c_{12}^- = v_- (c_{22}^- - 2c_{11}^-) + q_{11}^- (q_{11}^- - q_{22}^-)/2 = \\ \qquad = \frac{2}{3} \left[g_2^- (f_1^- + h_3^- + 3g_2^-) - 2h_2^- (f_1^- + h_3^- + 5g_2^-) v_- \right] + 4(h_2^- v_-)^2, \\ \clubsuit \quad c_{13}^- = \frac{2}{9} \left\{ 3(g_1^- - f_1^- g_3^- + g_2^- g_3^-) - \left[2(f_1^- + h_3^-)^2 + 3(h_1^- + h_2^- g_3^- - f_1^- h_3^-) + \right. \\ \left. + 6g_2^- (f_1^- + g_2^- + h_3^-) \right] + 5h_2^- (f_1^- + 3g_2^- + h_3^-) v_-^2 - 8(h_2^-)^2 v_-^3 \right\}, \\ \clubsuit \quad c_{23}^- = -v_- c_{22}^- - (q_{22}^-)^2 = -\frac{2}{9} \left[2(f_1^- + h_3^-)^2 + h_2^- (f_1^- + 3g_2^- + h_3^-) v_- - 4(h_2^- v_-)^2 \right]. \end{array}$$

Examples & Applications

A multi-parametric example



A multi-parametric example



Bifurcation set to be completed. See the preprint by A.Algaba, E. Freire, E. Gamero and C. García, Bifurcation analysis of planar nilpotent reversible systems.

Electronic converters

- Electronic Power converters are Switching-Mode Power Supplies (SMPS).
- Basically, they are built using semiconductor switches (diodes, transistors) and energy storage elements (inductors, capacitors)
- Examples: Rectifiers AC-DC, inverters DC-AC, DC-DC converters (buck, boost, buck-boost)
- Their mathematical models commonly lead to nonsmooth dynamical systems

DC-DC converters



The BOOST converter





• The value of q_1 stands for a controlled switch

• when q_1 is turned ON \rightarrow current in L increases and energy is stored in it

• when q_1 is turned OFF \rightarrow the stored energy in L is dropped and the polarity of the L voltage changes so that it adds to the input voltage

The goal is to get $V_{out} > \mathsf{E}$

The BOOST converter



The model of the system is given by

$$L\frac{di_L}{dt} = V_{in} - r_L i_L - u \cdot v_c \qquad (V_{in} = \mathsf{E})$$
$$C\frac{dv_c}{dt} = u \cdot i_L - \frac{v_c}{R}$$

where $v_c > 0$ is the capacitor voltage, $i_L > 0$ is the inductor current and $u \in \{0,1\}$ is the control action. Input voltage is assigned as V_{in} , r_L is the equivalent series resistance of the inductor, R is the resistive load, C and L are the capacitor and inductor, respectively.

The BOOST converter

To analyze the model, the system is normalized

$$(i_L, v_c) = \left(V_{in} \sqrt{\frac{C}{L}} x, V_{in} y \right)$$
 and $t = \tau \sqrt{LC};$

and new parameters are taken, namely

$$b = r_L \sqrt{\frac{C}{L}}$$
 and $a = \frac{1}{R} \sqrt{\frac{L}{C}};$

so that the boost converter model in dimensionless normal form is

$$\dot{x} = 1 - bx - u \cdot y$$
$$\dot{y} = u \cdot x - ay,$$

for x, y > 0, $u = \{0, 1\}, b \ge 0$, and a > 0.

The strategy at the BOOST converter







 $E_0 \begin{cases} \dot{x} = 1 - b \, x - y \\ \dot{y} = x - a \, y \end{cases}$





-Frequency variable Control -Frequency constant Control using PWM (Pulse-Width Modulation)

We want to regulate the output voltage to a desired value $y = v_c/V_{in} = y_r > 1$, ensuring robustness under parameter variation of a, produced by load changes in R.

A washout filter is used: the inductor current x can be filtered to get a new signal x_F by means of a washout filter given by the transfer function

$$G_F(s) = \frac{X_F(s)}{X(s)} = \frac{s}{s+w} = 1 - \frac{w}{s+w},$$

where w is the reciprocal of the filter constant and x_F is the filter output.

A differential equation is added, $\dot{z} = w(x - z)$, where z is a new state satisfying the output equation $x_F = x - z$.



The SMC strategy consists in the choice of

$$\Sigma = \{ \mathbf{x} = (x, y, z) \in \mathbb{R}^3 : h(\mathbf{x}) = y - y_r + k(x - z) = 0 \}$$

as the switching boundary, where we want to be located the pseudo-equilibrium point.

Control parameter

We use the vector fields defined by

$$\mathbf{F}^{+}(\mathbf{x}) = \begin{bmatrix} 1 - bx - y \\ x - ay \\ w(x - z) \end{bmatrix} \text{ and } \mathbf{F}^{-}(\mathbf{x}) = \begin{bmatrix} 1 - bx \\ -ay \\ w(x - z) \end{bmatrix},$$

for $h(\mathbf{x}) > 0$ and $h(\mathbf{x}) < 0$, respectively.

In what follows, we assume for simplicity b = 0 (ideal inductance), w = 1 and $y_r = 2$ (doubling the voltage).

Therefore, we work with

$$\begin{split} \Sigma &= \{ \mathbf{x} = (x, y, z) \in \mathbb{R}^3 : h(\mathbf{x}) = y - 2 + k(x - z) = 0 \} \\ \text{and} \\ \mathbf{F}^+(\mathbf{x}) &= \begin{bmatrix} 1 - y \\ x - ay \\ x - z \end{bmatrix} \text{ and } \mathbf{F}^-(\mathbf{x}) = \begin{bmatrix} 1 \\ -ay \\ x - z \end{bmatrix}, \begin{array}{c} \text{Control parameter} \\ \mathbf{x} - z \end{bmatrix}, \\ \text{for } h(\mathbf{x}) > 0 \text{ and } h(\mathbf{x}) < 0, \text{ respectively.} \end{split}$$

We need to compute $\nabla h \cdot \mathbf{F}^{\pm}$ on Σ , and look for Σ_{as} , namely

$$\nabla h \cdot \mathbf{F}^{+}(\mathbf{x})\big|_{\Sigma} = [k, 1, -k] \begin{bmatrix} 1-y\\ x-ay\\ x-z \end{bmatrix} = k(1-y) + x - ay - k(x-z) = x + (1-a-k)y + k - 2 < 0,$$

and

$$\nabla h \cdot \mathbf{F}^{-}(\mathbf{x})\big|_{\Sigma} = [k, 1, -k] \begin{bmatrix} 1\\ -ay\\ x-z \end{bmatrix} = k - ay - k(x-z) = (1-a)y + k - 2 > 0.$$

The tangency lines are

$$T_{-} = \{ \mathbf{x} \in \Sigma : (1-a)y + k - 2 = 0 \},\$$

$$T_{+} = \{ \mathbf{x} \in \Sigma : x + (1-a-k)y + k - 2 = 0 \}.$$

The double tangency point occurs where the tangency lines intersect transversally, i.e., at the point $\hat{\mathbf{x}} = (k\hat{y}, \hat{y}, \hat{z})$, with

where we assume $a \neq 1$.

To have $\hat{y} > 0$, we will consider $\operatorname{sign}(k-2) = \operatorname{sign}(a-1)$.



The pseudo-equilibrium point is at $(\tilde{x}, \tilde{y}, \tilde{z}) = (4a, 2, 4a)$ to be in Σ_{as} only if k > 2a. Thus, we must expect to have the compound bifurcation for k = 2a.

First, we need to check the conditions (H2) to have a TS-point. Computations lead to the two conditions

$$k > 2, \quad 1 < a < \frac{k - 3 + \sqrt{17 - 34k + 37k^2 - 20k^3 + 4k^4}}{2(k - 2)}$$

Assuming these inequalities, we put the system in the canonical form and compute the critical coefficients.

We obtain for the new coefficients

$$v_{-} = \frac{1 - a^{2}(k - 2) - a(3 - 3k + k^{2})}{(a - 1)\omega_{+}\omega_{-}}, \quad v_{+} = \frac{(a - k)(k - 2)}{\omega_{+}\omega_{-}},$$

along with

$$\begin{split} g_2^- &= h_1^- = h_2^- = 0, \qquad \qquad g_1^- = -\frac{1}{k\omega_-\omega_+}, \qquad f_1^- = -\frac{1}{\omega_-}, \\ g_3^- &= \frac{a^2k + a(1-k+k^2) - 1}{(a-1)k\omega_+}, \qquad h_3^- = -\frac{a}{\omega_-}, \qquad \qquad f_1^+ = -\frac{k-1}{k\omega_+}, \\ g_1^+ &= \frac{k^2 + (a-2)k + 1}{k^2\omega_+^2}, \qquad \qquad g_2^+ = -\frac{k^2 + (a-1)k + 1}{k\omega_+}, \qquad h_1^+ = -\frac{a-1}{k\omega_-\omega_+}, \\ g_3^+ &= -\frac{(k^2 + (a-2)k + 2 - a)\omega_-}{(a-1)\omega_+^2}, \qquad h_2^+ = \frac{a-1}{\omega_-}, \qquad \qquad h_3^+ = \frac{k-1}{\omega_+}, \end{split}$$

where the two conditions

$$\omega_{-}^{2} = a(k-2) > 0, \quad \omega_{+}^{2} = \frac{(k-1)^{3} + (k-3)a - (k-2)a^{2}}{a-1} > 0,$$

are guaranteed from Hypothesis (H2).

The bifurcation curve in the (a, k)-plane of parameters turns out to be

$$v_{-}v_{+} - 1 = (2a - k)\hat{y} = 0,$$

i.e.,

$$k = 2a.$$

Moreover, we have

$$v_+|_{k=2a} = \frac{2a(1-a)}{\omega_+\omega_-} < 0$$
 and $v_-|_{k=2a} = -\frac{(1-a)^2 + 5a^2}{\omega_+\omega_-} < 0$

for all a > 1. The sliding bifurcation is supercritical, since

$$\kappa_S \Big|_{k=2a} = \frac{2}{\left((1-a)^2 + 5a^2\right)^{\frac{3}{2}}} > 0.$$

Regarding the crossing bifurcation, we obtain

$$\sigma\big|_{k=2a} = \frac{2(5 - 4a + 16a^2 + 24a^3)}{3(1 - 2a + 6a^2)^{\frac{3}{2}}} > 0,$$

and

$$\kappa_C = \frac{2}{3((1-a)^2 + 5a^2)^2} > 0,$$

so that the bifurcating periodic orbit for k < 2a is of **stable node** type.
Simulation results



3.3 = k > 2a = 3

Simulation results



k = 2a = 3

Simulation results



2.5 = k < 2a = 3

Undesirable oscillation observed in laboratory due to the TS-point in the DC-DC Boost converter



Conclusions

- The TS compound bifurcation has been characterized for piecewise linear systems, and a procedure for computing the essential coefficients has been provided.
- Several examples have been shown, and in particular the bifurcation is detected in DC-DC converters under SMC strategy with a washout filter.
- A codimension-two unfolding is still needed, gaining information on secondary bifurcation curves.
- A question to solve: Can this study with linear pieces serve as a normal form for nonlinear vector fields?

Thank you for your attention!